

THERMOELASTIC WAVES IN AN INFINITE SOLID CAUSED BY A LINE HEAT SOURCE

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ABSTRACT. The generalized thermoelasticity theory recently developed by Green and Naghdi is employed to investigate thermoelastic interactions caused by a continuous line heat source in a homogeneous isotropic unbounded solid. Hankel–Laplace transform technique is used to solve the problem. Explicit expressions, for stress and temperature fields, are obtained for small time approximation. Numerical values are displayed graphically. Our results show that this theory predicts an infinite speed for heat propagation in general, and includes the second sound phenomena as a special case.

KEY WORDS AND PHRASES. Generalized thermoelasticity, thermoelastic waves.

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1. INTRODUCTION.

In recent years many papers have been devoted to the development of the generalized theory of thermoelasticity, which predicts a finite speed for heat propagation. Lord and Shulman [1], based on a modified Fourier's law, developed a generalized theory of thermoelasticity whose governing system of equations are entirely hyperbolic and hence predict finite speed for heat propagation. Green and Lindsay [2], based on an entropy production inequality proposed by Green and Laws [3], developed a temperature–rate dependent thermoelasticity that includes the temperature–rate among constitutive variables and also predicts a finite speed for heat propagation. The applications of these theories have been examined extensively by many authors (see [4]–[9]).

Recently Green and Naghdi [10] re-examined the basic postulates of thermomechanics and postulated three type of constitutive response functions for the thermal phenomena. The nature of these three types of constitutive equations is [11] such that when the respective theories are linearized, type I theory is the same as the classical heat conduction theory (based on Fourier's law); type II theory predicts a finite speed for heat propagation and involves no energy dissipation; type III theory permits propagation of thermal signals at both infinite and finite speeds [12] and there is structural difference between these field equations and those developed in [1] and [2].

The aim of the present paper is to study the thermoelastic interactions caused by a continuous line heat source in a homogeneous and isotropic infinite solid by employing the above mentioned type III theory. We use Hankel-Laplace transform to find the explicit expressions for stress and temperature fields in small time intervals. Our results indicate that this theory generally predicts diffusion type of heat propagation and includes the wave type of heat propagation as a special case. The counterparts of this problem in the context of theories developed in [1] and [2] have been studied respectively by Sherief and Anwar [8] and Chandrasekharaiah and Murthy [9]. Due to the structural difference between these theories, our results differ from those obtained in [8] and [9].

2. PRELIMINARIES

We consider a homogeneous isotropic elastic solid occupying the whole space. The governing system of equations in thermoelasticity of type III developed in [12] are

$$(\lambda + \mu)u_{j,ij} + \mu u_{i,jj} - \gamma\theta_{,i} + \rho f_i = \rho \bar{u}_i, \quad (2.1)$$

$$\rho c \bar{\theta} + \gamma \theta \cdot \bar{u}_{i,i} = \rho \dot{Q} + k \dot{\theta}_{,ii} + k^* \theta_{,ii}, \quad (2.2)$$

$$\sigma_{ij} = \lambda u_{k,k} \delta_{ij} + \mu (u_{i,j} + u_{j,i}) - \gamma \delta_{ij} \theta, \quad (2.3)$$

where λ and μ are Lamé constants, $\gamma = \frac{1}{3} E \beta^* / (1 - 2\nu)$, E is Young's modulus, ν is the Poisson ratio, β^* is the coefficient of volume expansion, u_i are the components of the displacement vector, ρ is the mass density, θ is the temperature deviation above the initial temperature θ_0 , σ_{ij} are the components of stress tensor, k is the thermal conductivity, k^* is a constant, c is the specific heat for processes with invariant strain tensor, Q and f_i are, respectively, the heat source and the components of the body force, measured per unit volume.

In above equations, the notation of Cartesian tensor is employed, superposed dots denote the time derivatives and a comma followed by the index i denotes the partial derivative with respect to x_i . Using the nondimensional variables

$$\begin{aligned} x'_i &= x_i/l, & t' &= ta_0/l, & \theta' &= \theta/\theta_0, & u'_i &= u_i/l, \\ \sigma'_{ij} &= \sigma_{ij}/\mu, & \rho' &= \rho/\rho_0, & f'_i &= f_i l/a^2, & Q' &= Q l/a^3, \end{aligned}$$

where l is a standard length, a_0 is a standard speed, ρ_0 is a standard mass density, the basic equations (2.1)–(2.3) reduce to the following (dropping primes for convenience):

$$\rho \alpha_1 \bar{u}_i = \alpha_2 u_{j,ij} + u_{i,jj} - \alpha_3 \theta_{,i} + \rho \alpha_4 f_i, \quad (2.4)$$

$$\theta_{,ii} + \alpha_4 \dot{\theta}_{,ii} + \rho \alpha_5 \dot{Q} = \rho \alpha_6 \bar{\theta} + \alpha_7 \bar{u}_{i,i}, \quad (2.5)$$

$$\sigma_{ij} = \frac{\lambda}{\mu} u_{k,k} \delta_{ij} + (u_{i,j} + u_{j,i}) - \alpha_3 \theta, \quad (2.6)$$

with

$$\alpha_1 = \rho \cdot a^2 / \mu, \quad \alpha_2 = (\lambda + \mu) / \mu, \quad \alpha_3 = \gamma \theta_0 / \mu, \quad \alpha_4 = ka_0 / (k^* t),$$

$$\alpha_5 = \rho \cdot a^4 / (k^* \theta_0), \quad \alpha_6 = \rho \cdot ca^2 / k^*, \quad \alpha_7 = \gamma a^2 / k^*.$$

3. SOLUTION IN LAPLACE TRANSFORM DOMAIN

In the present paper we consider an infinite solid containing a line heat source situated along the x_3 -axis, that is,

$$f_i = 0 \quad \text{and} \quad Q = \frac{1}{2\pi r} Q_0 \delta(r) H(t), \quad (3.1)$$

where $\delta(t)$ is the Dirac delta function, $H(t)$ is the Heaviside unit step function, $r = (x_1^2 + x_2^2)^{1/2}$ and Q_0 is a constant. The resulting thermoelastic interactions are axisymmetric in nature so that the displacement vector has only the radial component $u = u(r, t)$ and the stress tensor has only two components σ_r and σ_ϕ which are normal stresses in the radial and transverse directions.

In the context of the problem considered, the regularity conditions are taken as

$$(\sigma_r, \sigma_\phi, \theta) \longrightarrow (0, 0, 0) \quad \text{as} \quad r \longrightarrow \infty, \quad (3.2)$$

and the initial conditions at $t = 0$ are

$$u = \dot{u} = \theta = \dot{\theta} = 0. \quad (3.3)$$

Transforming equations (2.4)–(2.6) into cylindrical coordinates, with $f_i = 0$, we obtain

$$\rho \alpha_1 \frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial r} \left[(\alpha_2 + 1) \left(\frac{\partial u}{\partial r} + \frac{u}{r} \right) - \alpha_3 \theta \right], \quad (3.4)$$

$$\theta_{,ii} + \alpha_4 \dot{\theta}_{,ii} + \rho \alpha_5 \dot{Q} = \rho \alpha_6 \dot{\theta} + \alpha_7 \frac{1}{r} \frac{\partial}{\partial r} (r \partial^2 u / \partial t^2), \quad (3.5)$$

$$\sigma_r = (\alpha_2 + 1) \frac{\partial u}{\partial r} + (\alpha_2 - 1) \frac{u}{r} - \alpha_3 \theta, \quad (3.6)$$

$$\sigma_\phi = (\alpha_2 - 1) \frac{\partial u}{\partial r} + (\alpha_2 + 1) \frac{u}{r} - \alpha_3 \theta. \quad (3.7)$$

Introducing the thermoelastic potential function φ defined by

$$u = \frac{\partial \varphi}{\partial r}, \quad (3.8)$$

equations (3.4) and (3.5) reduce to

$$\rho \alpha_1 \frac{\partial^2 \varphi}{\partial t^2} = (\alpha_2 + 1) \nabla^2 \varphi - \alpha_3 \theta, \quad (3.9)$$

$$\nabla^2 \theta + \frac{\partial}{\partial t} [\alpha_4 \nabla^2 \theta + \rho \alpha_5 Q] = \frac{\partial^2}{\partial t^2} [\rho \alpha_6 \theta + \alpha_7 \nabla^2 \varphi], \quad (3.10)$$

where $\nabla^2 = \frac{\partial}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r}$. Using equations (3.9) and (3.10), equations (3.6) and (3.7) reduce to

$$\sigma_r = \rho \alpha_1 \frac{\partial^2 \varphi}{\partial t^2} - \frac{2}{r} \frac{\partial \varphi}{\partial r}, \quad (3.11)$$

$$\sigma_\phi = \rho \alpha_1 \frac{\partial^2 \varphi}{\partial t^2} - 2 \frac{\partial^2 u}{\partial r^2}. \quad (3.12)$$

Elimination of θ between equations (3.9) and (3.10) gives

$$\nabla^4 \varphi + \alpha_4 \frac{\partial}{\partial t} (\nabla^4 \varphi) - b_1 \frac{\partial^3}{\partial t^3} (\nabla^2 \varphi) - b_2 \frac{\partial^2}{\partial t^2} (\nabla^2 \varphi) + b_3 \frac{\partial^4 \varphi}{\partial t^4} + b_4 \frac{\partial Q}{\partial t} = 0, \quad (3.13)$$

where

$$b_1 = \rho \alpha_1 \alpha_4 / (\alpha_2 + 1), \quad b_2 = \rho \alpha_6 + (\rho \alpha_1 + \alpha_3 \alpha_7) / (\alpha_2 + 1),$$

$$b_3 = \rho^2 \alpha_1 \alpha_6 / (\alpha_2 + 1), \quad b_4 = \rho \alpha_3 \alpha_5 / (\alpha_2 + 1).$$

Applying the Laplace transform, defined by

$$\bar{g}(r, p) = \int_0^\infty g(r, t) \exp\{-pt\} dt, \quad \operatorname{Re}(p) > 0,$$

to equations (3.11)–(3.13) under the homogeneous initial conditions, with Q given by (7), we find that

$$\bar{\sigma}_r = (b_5^2 p^2 - \frac{2}{r} \frac{d}{dr}) \bar{\varphi}, \quad (3.14)$$

$$\bar{\sigma}_\phi = (b_5^2 p^2 - 2 \frac{d}{dr^2}) \bar{\varphi}, \quad (3.15)$$

$$[(1 + \alpha_4 p) \nabla^4 - (b_1 p + b_2) p^2 \nabla^2 + b_3 p^4] \bar{\varphi} + \frac{A}{r} \delta(r) = 0, \quad (3.16)$$

where $A = \frac{1}{2} b_4 Q_0 / \pi$ and $b_5^2 = \rho \alpha_1$. We may rewrite equation (3.16) in the form

$$(\nabla^2 - \lambda_1^2)(\nabla^2 - \lambda_2^2) \bar{\varphi} = - \frac{A \delta(r)}{r(1 + \alpha_4 p)}, \quad (3.17)$$

where λ_1^2 and λ_2^2 are roots of the characteristic equation

$$z^2 - \frac{b_1 p + b_2}{1 + \alpha_4 p} p^2 z + \frac{1}{1 + \alpha_4 p} b_3 p^4 = 0. \tag{3.18}$$

Applying the Hankel transform defined by

$$\hat{g}(\xi, p) = \int_0^\infty r J_0(\xi r) \bar{g}(r, p) dr,$$

where J_0 is the Bessel function of the first kind and of zero order, to equation (23), we find that

$$(\xi^2 + \lambda_1^2)(\xi^2 + \lambda_2^2)\hat{\varphi}(\xi, p) = -\frac{A}{1 + \alpha_4 p}. \tag{3.19}$$

Using the inverse Hankel transform defined by

$$\bar{g}(r, p) = \int_0^\infty \xi J_0(\xi r) \hat{g}(\xi, p) d\xi,$$

in equation (3.19), we obtain

$$\bar{\varphi}(r, p) = \frac{A}{1 + \alpha_4 p} [K_0(\lambda_1 r) - K_0(\lambda_2 r)] / (\lambda_1^2 - \lambda_2^2), \tag{3.20}$$

where K_0 is the modified Bessel function of the second kind and of zero order.

Using the equation (3.20) and the following recurrence relations of the modified Bessel functions of the second kind

$$\frac{\partial}{\partial r} K_0(r) = -K_1(r), \quad \text{and} \quad \frac{\partial}{\partial r} [r K_1(r)] = -r K_0(r),$$

where $K_1(r)$ is the modified Bessel function of second kind and of order one, the equations (3.14) and (3.15) become

$$\bar{\sigma}_r = \frac{A}{1 + \alpha_4 p} \sum_{i=1}^2 (-1)^{i-1} \left[b_3^2 p^2 K_0(\lambda_i r) + \frac{2}{r} \lambda_i K_1(\lambda_i r) \right] / (\lambda_1^2 - \lambda_2^2), \tag{3.21}$$

$$\bar{\sigma}_\phi = \frac{A}{1 + \alpha_4 p} \sum_{i=1}^2 (-1)^{i-1} \left[(b_5^2 p^2 - 2\lambda_i^2) K_0(\lambda_i r) - \frac{2}{r} \lambda_i K_1(\lambda_i r) \right] / (\lambda_1^2 - \lambda_2^2). \tag{3.22}$$

Taking the Laplace transform of both sides of equation (3.9) and using the equation (26), we find that

$$\bar{\theta} = \frac{A}{(1 + \alpha_4 p) \alpha_3} \left[(\lambda_1^2 - b_5^2 p^2) K_0(\lambda_1 r) - (\lambda_2^2 - b_5^2 p^2) K_0(\lambda_2 r) \right] / (\lambda_1^2 - \lambda_2^2). \quad (3.23)$$

The system of equations (3.21)–(3.23) gives rise to the solutions of stress and temperature fields in Laplace transform domain. The solutions in (r, t) domain can be obtained by inverting the Laplace transform.

4. SMALL-TIME SOLUTION IN (r, t) DOMAIN

It is a formidable task to find the inverse Laplace transform of equations (3.21)–(3.23). For this reason we have resorted to the case of small-time approximation. From equation (3.18), we find that

$$\lambda_{1,2} = p \left[\frac{(b_1 p + b_2) \pm \sqrt{(b_1 p + b_2)^2 - 4b_3(1 + \alpha_4 p)}}{2(1 + \alpha_4 p)} \right]^{1/2}. \quad (4.1)$$

For large p , expanding the above equation binomially in ascending powers of $1/p$ and retaining only necessary terms, we obtain

$$\lambda_1 \approx b_{10} p + b_{11}, \quad \lambda_2 \approx b_{20} p^{1/2}, \quad (4.2)$$

by neglecting terms of $O(p^{-2})$, where

$$b_{10} = (b_1/\alpha_4)^{1/2} = b_5/(\alpha_2+1)^{1/2}, \quad b_{20} = (b_3)^{1/2}, \\ b_{11} = \frac{1}{2}(b_1 b_2 \alpha_4 - b_3 \alpha_4^2 - b_1^2)/(\alpha_4)^{3/2}.$$

Similarly, we have

$$\frac{1}{(\lambda_1^2 - \lambda_2^2)(1 + \alpha_4 p)} = \frac{1}{b_1} \frac{1}{p^3} + b_0 \frac{1}{p^4} + O(p^{-5}), \quad (4.3)$$

where

$$b_0 = (2b_3 \alpha_4 - b_1 b_2)/b_1^2.$$

Substituting from equations (4.2) and (4.3) into equations (3.21)–(3.23), we obtain

$$\bar{\sigma}_r/A \approx b_5^2 \left[\frac{1}{b_1} p^{-1} + b_0 p^{-2} \right] K_0[(b_{10} p + b_{11})r] + \frac{2}{r} \left[(b_{10}/b_1) p^{-2} + (b_{11}/b_1 + b_{10} b_0) p^{-3} \right. \\ \left. + b_{11} b_0 p^{-4} \right] K_1[(b_{10} p + b_{11})r] - b_5^2 \left[\frac{1}{b_1} p^{-1} + b_0 p^{-2} \right] K_0[b_{20} p^{1/2} r] - \frac{2}{r} b_{20} \left[\frac{1}{b_1} p^{-5/2} \right. \\ \left. + b_0 p^{-7/2} \right] K_1[b_{20} p^{1/2} r], \quad (4.4)$$

$$\begin{aligned} \bar{\sigma}_\phi/A \approx & \left[\frac{b_5^2 - 2b_{10}^2}{b_1} p^{-1} + (b_5^2 b_0 - 2b_{10}^2 b_0 - 4b_{10} b_{11}/b_1) p^{-2} - 2(b_{11}^2/b_1 + 2b_{10} b_0) p^{-3} \right. \\ & \left. - 2b_{11}^2 b_0 p^{-4} \right] K_0[(b_{10} p + b_{11}) r] - \frac{2}{r} \left[(b_{10}/b_1) p^{-2} + (b_{11}/b_1 + b_{10} b_0) p^{-3} + b_{11} b_0 p^{-4} \right] \\ & K_{11}[(b_{10} p + b_{11}) r] - \left[(b_5^2/b_1) p^{-1} + (b_5^2 b_0 - 2b_{20}^2/b_1) p^{-2} - 2b_{20}^2 b_0 p^{-3} \right] K_0[b_{20} p^{1/2} r] \\ & + \frac{2}{r} b_{20} \left(\frac{1}{b_1} p^{-5/2} + b_0 p^{-7/2} \right) K_{11}[b_{20} p^{1/2} r], \end{aligned} \tag{4.5}$$

$$\begin{aligned} \alpha_3 \bar{\theta}/A \approx & \left[\frac{b_{10}^2 - b_5^2}{b_1} p^{-1} + (b_{10}^2 b_0 - b_5^2 b_0 + 2b_{10} b_{11}/b_1) p^{-2} + (b_{11}^2/b_1 + 2b_{10} b_{11} b_0) p^{-3} \right. \\ & \left. + b_{11}^2 b_0 p^{-4} \right] K_0[(b_{10} p + b_{11}) r] + \left[(b_5^2/b_1) p^{-1} + (b_5^2 b_0 - b_{20}^2/b_1) p^{-2} - b_{20}^2 b_0 p^{-3} \right] K_0[b_{20} p^{1/2} r]. \end{aligned} \tag{4.6}$$

From [13], we find the following inverse Laplace transforms

$$\mathcal{L}^{-1}\{p^{-n}\} = t^{n-1}/\Gamma(n), \tag{4.7}$$

$$\mathcal{L}^{-1}\{K_0(\alpha p)\} = H(t-\alpha)(t^2-\alpha^2)^{-1/2}, \tag{4.8}$$

$$\mathcal{L}^{-1}\{p^{-1}K_0(\alpha p)\} = H(t-\alpha)\cosh^{-1}(t/\alpha), \tag{4.9}$$

where $\Gamma(n)$ is the Gamma function. Using the shift property of the Laplace transform and the expansion technique for large p , we obtain

$$\begin{aligned} \mathcal{L}^{-1}\{p^{-1}K_0[(b_{10} p + b_{11}) r]\} &= \exp\{-mt\} \mathcal{L}^{-1}\left\{\frac{1}{p-m} K_0(b_{10} p r)\right\} \\ &= \exp\{-mt\} \mathcal{L}^{-1}\{(p^{-1} + mp^{-2} + m^2 p^{-3}) K_0(b_{10} p r)\}, \end{aligned} \tag{4.10}$$

where $m = b_{11}/b_{10}$. Using the convolution theorem of Laplace transform and the formulas (4.7) and (4.8), we get

$$\begin{aligned} \mathcal{L}^{-1}\{p^{-2}K_0(b_{10} p r)\} &= \mathcal{L}^{-1}\{p^{-2}\} * \mathcal{L}^{-1}\{K_0(b_{10} p r)\} \\ &= [t/\Gamma(2)] * [H(t-b_{10} r)(t^2-b_{10}^2 r^2)^{-1/2}] \\ &= \int_0^t (t-s) H(s-b_{10} r)(s^2-b_{10}^2 r^2)^{-1/2} ds \\ &= H(t-b_{10} r) \left[t \cosh^{-1}[t/(b_{10} r)] - (t^2-b_{10}^2 r^2)^{1/2} \right]. \end{aligned} \tag{4.11}$$

Similarly we obtain

$$\mathcal{L}^{-1}\{p^{-3}K_0(b_{10} p r)\} = H(t-b_{10} r) \left[\frac{1}{2} t^2 + \frac{1}{4} b_{10}^2 r^2 \cosh^{-1}[t/(b_{10} r)] - \frac{3}{4} t(t^2-b_{10}^2 r^2)^{1/2} \right]. \tag{4.12}$$

Substituting from equations (4.9), (4.11) and (4.12) into equation (4.10), we obtain

$$\mathcal{L}^{-1}\{p^{-1}K_0[(b_{10}p+b_{11})r]\} = H(t-b_{10}r)\exp\{-mt\}\left[(1+mt+\frac{1}{2}m^2t^2+\frac{1}{4}m^2b_{10}^2r^2) \times \right. \\ \left. \cosh^{-1}\left(\frac{t}{b_{10}r}\right) - m\left(1+\frac{3}{4}mt\right)(t^2-b_{10}^2r^2)^{1/2}\right], \quad (4.13)$$

Applying similiar procedure to other terms and making use of the following identities

$$\mathcal{L}^{-1}\{K_1(\alpha p)\} = \frac{t}{\alpha}H(t-\alpha)(t^2-\alpha^2)^{-1/2}, \quad (4.14)$$

$$\mathcal{L}^{-1}\{p^{-1}K_1(\alpha p)\} = \frac{1}{\alpha}H(t-\alpha)(t^2-\alpha^2)^{1/2}, \quad (4.15)$$

$$\mathcal{L}^{-1}\{K_0(\sqrt{\alpha p})\} = \frac{1}{2t}\exp\{-\frac{\alpha}{4t}\}, \quad (4.16)$$

$$\mathcal{L}^{-1}\{p^{-1/2}K_1(\sqrt{\alpha p})\} = \alpha^{-1/2}\exp\{-\frac{\alpha}{4t}\}, \quad (4.17)$$

we find that for small values of time

$$\sigma_r/A = H(t-b_{10}r)\exp\{-mt\}\left[\sigma_{r1}\cosh^{-1}\left(\frac{t}{b_{10}r}\right) + \sigma_{r2}(t^2-b_{10}^2r^2)^{1/2}\right] + \sigma_{r3}E_1\{b_{20}^2r^2/(4t)\} \\ + \sigma_{r4}\exp\{-b_{20}^2r^2/(4t)\}, \quad (4.18)$$

$$\sigma_\phi/A = H(t-b_{10}r)\exp\{-mt\}\left[\sigma_{\phi1}\cosh^{-1}\left(\frac{t}{b_{10}r}\right) + \sigma_{\phi2}(t^2-b_{10}^2r^2)^{1/2}\right] + \sigma_{\phi3}E_1\{b_{20}^2r^2/(4t)\} \\ + \sigma_{\phi4}\exp\{-b_{20}^2r^2/(4t)\}, \quad (4.19)$$

$$\alpha_3\theta/A = H(t-b_{10}r)\exp\{-mt\}\left[\theta_1\cosh^{-1}\left(\frac{t}{b_{10}r}\right) + \theta_2(t^2-b_{10}^2r^2)^{1/2}\right] + \theta_3E_1\{b_{20}^2r^2/(4t)\} \\ + \theta_4\exp\{-b_{20}^2r^2/(4t)\}, \quad (4.20)$$

where $E_1\{x\} = \int_x^\infty \frac{1}{s}\exp\{-s\}ds$, $x > 0$, is the exponential integral and

$$\sigma_{r1} = \frac{1}{8b_1}\left[8b_5^2 + 4b_5^2b_{11}^2r^2 + 4b_0b_1b_5^2b_{10}b_{11}r^2 + 3b_{10}b_{11} - 4b_1b_{10}^2 + b_0b_1b_{10}^2 \right. \\ \left. + b_0b_1b_{10}^3b_{11}r^2\right] + \frac{t}{16b_1}\left[16b_5^2m + 16b_0b_1b_5^2 - 32b_{11}^2 - 16b_{10}b_{11} - 16b_0b_1b_{10}^2 \right. \\ \left. - b_0b_1b_{10}b_{11}\right] + \frac{t^2}{2b_1}\left[m^2b_5^2 + 2mb_0b_1b_5^2 - b_0b_1b_{10}b_{11}\right], \quad (4.21)$$

$$\sigma_{r2} = \frac{1}{2b_1}\left[b_0b_1b_{10}^2 + 3b_{10}b_{11} - 2mb_5^2 - 2b_0b_1b_5^2\right] + \frac{t}{8b_1}\left[\frac{1}{r^2} + 2\frac{m}{r^2} + 7b_0b_1b_{10}b_{11} \right. \\ \left. - 12mb_1b_5^2 - 12mb_0b_1b_5^2\right] + \frac{t^2}{16b_{10}b_1}\left[\frac{8}{r^2} - \frac{1}{r}b_0b_1\right] + \frac{t^3}{8b_{10}r^2}b_0b_{11}, \quad (4.22)$$

$$\sigma_{r3} = \frac{1}{96b_1}\left[b_0b_1b_{20}^6r^4 + 6b_{20}^4r^2 - 48b_5^2 - 12b_0b_1b_5^2b_{20}^2r^2\right] + \frac{t}{16b_1}\left[4b_0b_1b_{20}^2 + b_0b_1b_{20}^4r^2 \right. \\ \left. + 8b_{20}^2 - 8b_0b_1b_5^2\right], \quad (4.23)$$

$$\sigma_{r4} = \frac{t}{24b_1} \left[12b_0b_1b_5^2 - 6b_{20}^2 - b_0b_1b_{20}^4r^2 \right] - \frac{t^2}{24b_1} \left[24\frac{1}{r^2} + 5b_0b_1b_{20}^2 \right] - \frac{1}{3}b_0t^3/r^2, \quad (4.24)$$

$$\begin{aligned} \sigma_{\phi 1} = & \frac{1}{8b_1} \left[(8b_5^2 - 8b_{10}^2 - 2b_{10}b_{11} - b_0b_{11} - b_0^2b_1b_{10}) + 4(b_0b_1b_5^2 - 2b_0b_1b_{10}^2 - 4b_{10}b_{11})b_{10}b_{11}r \right. \\ & \left. + (4b_5^2b_{11}^2 - 12b_{10}^2b_{11}^2 - 8b_0b_1b_{10}^3 - b_0b_1b_{10}^3b_{11})r^2 \right] + \frac{t}{16b_1} \left[(16mb_5^2 + 16b_0b_1b_5^2 - 16b_0b_1b_{10}^2 \right. \\ & \left. - 48b_{10}b_{11} + b_0b_1b_{10}b_{11}) - 8b_0b_1b_{10}^2b_{11}^2r^2 \right] + \frac{t^2}{2b_1} \left[(m^2b_5^2 + 2b_0b_1mb_5^2 - 14b_{11}^2 - 4b_0b_1b_{10} \right. \\ & \left. - 3b_0b_1b_{10}b_{11}) - \frac{1}{3}b_0b_{11}^2t^3 \right], \quad (4.25) \end{aligned}$$

$$\begin{aligned} \sigma_{\phi 2} = & \frac{1}{18b_1} \left[(36b_{10}^2 + 27b_0b_1b_{10}^2 + 45b_{10}b_{11} - 18b_5^2 - 18b_0b_1b_5^2) + 16b_0b_{10}^2b_{11}^2r^2 \right] \\ & + \frac{t}{8b_1b_{10}} \left[(24b_0b_1b_{10}^2b_{11} + 24b_0b_1b_{10}^2 + 84b_{10}b_{11}^2 - 6m^2b_5^2b_{10} - 12b_0b_5^2b_{10}b_{11}) \right. \\ & \left. - 7b_0b_1b_{10}^2b_{11}r - (3b_{11} + b_0b_1b_{10} + 8b_{10})\frac{1}{r^2} \right] + t^2 \left[\frac{11}{18}b_0b_{11}^2 + \frac{1}{16b_1b_{10}}(b_0b_1b_{11} \right. \\ & \left. - 12b_{11} - 4b_0b_1b_{10})\frac{1}{r^2} \right] - \frac{1}{8}b_0b_{11}t^3/(b_{10}r^2), \quad (4.26) \end{aligned}$$

$$\begin{aligned} \sigma_{\phi 3} = & \frac{1}{96b_1} \left[-48b_5^2 + 18b_{20}^4r^2 - 12b_0b_1b_5^2b_{20}^2r^2 - b_0b_1b_{20}^6r^4 \right] \\ & + \frac{t}{16b_1} \left[8b_{20}^2 - 8b_0b_1b_5^2 - 4b_0b_{20}^2 - b_0b_1b_{20}^4r^2 \right], \quad (4.27) \end{aligned}$$

$$\sigma_{\phi 4} = t \left[\frac{1}{2}b_0b_5^2 - \frac{3}{4}b_{20}^2/b_1 + \frac{1}{24}b_0b_{20}^4r^2 \right] + t^2 \left[(b_1r^2)^{-1} + \frac{5}{24}b_0b_{20}^2 \right] + \frac{1}{3}b_0t^3/r^2, \quad (4.28)$$

$$\begin{aligned} \theta_1 = & \frac{1}{4b_1} \left[4(b_{10}^2 - b_5^2) + (7b_{10}^2b_{11}^2 + 10b_0b_{10}^3b_{11} - 2b_5^2b_{11}^2 - 2b_0b_1b_5^2b_{10}b_{11})r^2 \right] \\ & + \frac{t}{4b_1} \left[(12b_{10}b_{11} + 4b_0b_{10}^2 - 4\gamma b_5^2 - 4b_0b_1b_5^2) + b_0b_1b_{10}^2b_{11}^2r^2 \right] \\ & + \frac{t^2}{2b_1} \left[7b_{11}^2 + 4b_0b_1b_{10}b_{11} - 2mb_0b_1b_5^2 - 2m^2b_5^2 \right] + \frac{1}{6}b_0b_{11}^2t^3, \quad (4.29) \end{aligned}$$

$$\begin{aligned} \theta_2 = & (b_0b_5^2 + mb_5^2/b_1 - b_0b_{10}^2 - 3b_{10}b_{11}/b_1 - \frac{4}{9}b_0b_{10}^2b_{11}^2r^2) \\ & + t \left(\frac{3}{2}mb_0b_5^2 + \frac{3}{4}m^2b_5^2/b_1 - \frac{21}{4b_1}b_{11}^2 - 3b_0b_{10}b_{11} \right) - \frac{11}{36}b_0b_{11}^2t^2, \quad (4.30) \end{aligned}$$

$$\begin{aligned} \theta_3 = & \left(\frac{1}{2}b_5^2/b_1 + \frac{1}{8}b_0b_5^2b_{20}^2 - \frac{1}{8}b_{20}^4r^2/b_1 - \frac{1}{128}b_0b_{20}^6r^4 \right) \\ & + t \left(\frac{1}{2}b_0b_5^2 - \frac{1}{2}b_{20}^2/b_1 - \frac{1}{8}b_0b_{20}^4r^2 \right) - \frac{1}{2}b_0b_{20}^2t^2, \quad (4.31) \end{aligned}$$

$$\theta_4 = \frac{t}{2}(b_{20}^2/b_1 - b_0b_5^2 + \frac{r^2}{16}b_0b_{20}^4) + \frac{3}{8}b_0b_{20}^2t^2. \quad (4.32)$$

Due to the presence of the exponential integral and the exponent functions in the expressions of equations (4.18)–(4.20), it can be seen that the effect of input (7) is felt throughout the medium instantly and hence it shows that according to this theory heat travels at an infinite speed.

5. HEAT PROPAGATION WITH A FINITE SPEED

As we have pointed out earlier, this theory generally predicts a diffusion type of heat

propagation and includes the finite wave propagation of heat as a special case. In this section we consider this special case when $k^* \gg k$, that is, $\alpha_4 \approx 0$, $b_1 \approx 0$ and equation (4.1) becomes

$$\lambda_i = c_i p, \quad i = 1, 2, \quad (5.1)$$

where

$$c_i = \left[b_2 \pm (b_2^2 - 4b_3)^{1/2} / 2 \right]^{1/2}. \quad (5.2)$$

Using equation (5.1) in equations (3.21)–(3.23), we obtain

$$c_i \bar{\sigma}_r / A = \sum_{i=1}^2 (-1)^{i-1} \left[b_5^2 K_0(c_i p r) + \frac{2}{r} c_i p^{-1} K_1(c_i p r) \right], \quad (5.3)$$

$$c_i \bar{\sigma}_\phi / A = \sum_{i=1}^2 (-1)^{i-1} \left[(b_5^2 - 2c_i^2) K_0(c_i p r) - \frac{2}{r} c_i p^{-1} K_1(c_i p r) \right], \quad (5.4)$$

$$c_i \alpha_3 \bar{\theta} / A = \sum_{i=1}^2 (-1)^{i-1} (c_i^2 - b_5^2) K_0(c_i p r), \quad (5.5)$$

where $c_i = (b_2^2 - 4b_3)^{1/2}$.

Making use of identities (4.8) and (4.15), we find the solution for stress and temperature fields as

$$c_i \sigma_r / A = \sum_{i=1}^2 (-1)^{i-1} H(t - c_i r) \left[b_5^2 (t^2 - c_i^2 r^2)^{-1/2} + \frac{2}{r^2} (t^2 - c_i^2 r^2)^{1/2} \right], \quad (5.6)$$

$$c_i \sigma_\phi / A = \sum_{i=1}^2 (-1)^{i-1} H(t - c_i r) \left[(b_5^2 - 2c_i^2) (t^2 - c_i^2 r^2)^{-1/2} - \frac{2}{r^2} (t^2 - c_i^2 r^2)^{1/2} \right], \quad (5.7)$$

$$c_i \alpha_3 \theta / A = \sum_{i=1}^2 (-1)^{i-1} (c_i^2 - b_5^2) H(t - c_i r) (t^2 - c_i^2 r^2)^{-1/2}. \quad (5.8)$$

In this special case, solutions (5.6)–(5.8) indicate that thermal signals propagate at a finite speed since the Heaviside unit step function appears in all terms.

6. NUMERICAL RESULTS AND CONCLUSIONS

The numerical values of stress and temperature fields at time $t = 0.5$ have been calculated and displayed in figure 1–3 along the r -axis. To obtain these numerical values, we have taken that $b_2 = 5.2$, $b_3 = 3.35$, $b_4 = 2.25$, and $b_5^2 = 2.54$. In the general case, we assumed that $\alpha_4 = 3.1$ and $b_1 = 2.8$.

From Fig. 1 and 2, we note that in the general case the magnitudes of the radial stress and the tangential stress decrease from $-\infty$ at $r = 0$ to zero as r tends to infinity without

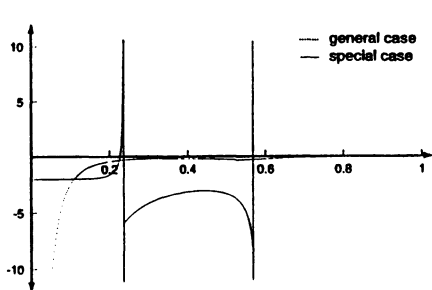


Fig.1. σ_r/A vs. r for $t = 0.5$

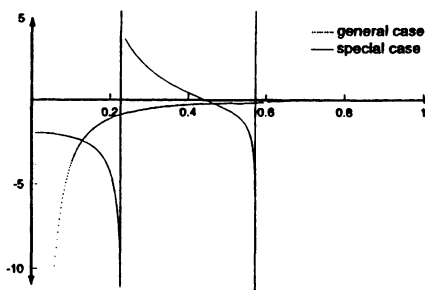


Fig.2. σ_ϕ/A vs. r for $t = 0.5$

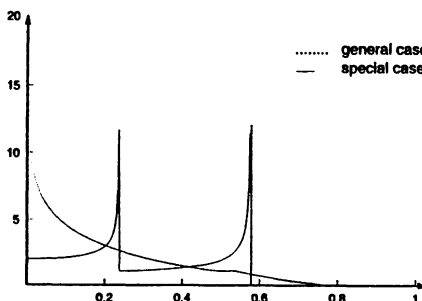


Fig.3. $\alpha_3\theta/A$ vs. r for $t = 0.5$

any jumps. However, in the undamped case, the corresponding magnitudes suffer two infinite jumps at $r = t/c_1 \approx 0.237113$ and $r = t/c_2 \approx 0.576053$. From Fig.3, we see that in the general case the temperature decreases from $+\infty$ at $r = 0$ to zero as r tends to infinity while the temperature in the undamped case suffers two infinite jumps at $r \approx 0.237113$ and $r \approx 0.576053$ and vanishes for $r > 0.576053$.

It is also apparent that, in general, this theory predicts a diffusion type of thermal propagation. The values of stress and temperature fields damp out gradually as r increases. In the special case (undamped case) when $k^* \gg k$, this theory predicts finite speed for heat propagation. In this special case, stress and temperature fields vanish identically for $r > t/c_2 \approx 0.576053$.

We also note that in the special case, both the stress and the temperature fields have finite values at $r = 0$, which is quite unusual for a continuous line heat source input given by (3.1). However, if we set $k = 0$ in the heat conduction equation (2.2), we obtain a hyperbolic heat equation which contains \dot{Q}

$$\rho c \ddot{\theta} + \gamma \theta_{,i,i} = \rho \dot{Q} + k^* \theta_{,ii}$$

In this case we get an impulsive heat source \dot{Q} when we take Q as a continuous heat source given by (3.1). According to solutions (5.6)–(5.8), the precise values for stress and temperature fields at $r = 0$ are

$$\sigma_r/A = -1/t, \quad \sigma_\phi/A = -1/t, \quad \theta/A = 1/t.$$

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